# <span id="page-0-0"></span>Fast Corotated Elastic SPH Solids with Implicit Zero-Energy Mode Control

TASSILO KUGELSTADT, RWTH Aachen University, Germany JAN BENDER, RWTH Aachen University, Germany JOSÉ ANTONIO FERNÁNDEZ-FERNÁNDEZ, RWTH Aachen University, Germany STEFAN RHYS JESKE, RWTH Aachen University, Germany FABIAN LÖSCHNER, RWTH Aachen University, Germany ANDREAS LONGVA, RWTH Aachen University, Germany

## ACM Reference Format:

Tassilo Kugelstadt, Jan Bender, José Antonio Fernández-Fernández, Stefan Rhys Jeske, Fabian Löschner, and Andreas Longva. 2021. Fast Corotated Elastic SPH Solids with Implicit Zero-Energy Mode Control. Proc. ACM Comput. Graph. Interact. Tech. 4, 3 (September 2021), [2](#page-1-0) pages. <https://doi.org/10.1145/3480142>

# 1 KERNEL GRADIENT CORRECTION

Computing the deformation gradient using the standard SPH gradient in Eq. (6) leads to artifacts since it is not  $1^\text{st}-$ order consistent. To solve this problem we analyze the error of Eq. (6) by replacing  $A(X_i)$  with its Taylor approximation at the point  $X_i$ 

$$
\sum_{j \in \mathcal{N}_i^0} V_j A_j \nabla W_{ij} = \sum_{j \in \mathcal{N}_i^0} V_j \left( A_i + \nabla A_i \cdot \mathbf{X}_{ji} + O(\mathbf{X}_{ji}^2) \right) \nabla W_{ij}
$$
  
= 
$$
\sum_{j \in \mathcal{N}_i^0} V_j A_i \nabla W_{ij} + \sum_{j \in \mathcal{N}_i^0} \left( V_j \nabla W_{ij} \otimes \mathbf{X}_{ji} \right) \nabla A_i + O(\mathbf{X}_{ji}^2),
$$

where  $X_{ji} = X_j - X_i$  and ⊗ denotes the dyadic product of two vectors  $\mathbf{a} \otimes \mathbf{b} = \mathbf{a} \mathbf{b}^T$ . To solve for  $\nabla A$ we subtract the first and third term and multiply with the inverse of the matrix from the second term. This yields a 1<sup>st</sup>-order consistent approximation of the gradient

$$
\nabla A_i \approx \sum_{j \in \mathcal{N}_i^0} V_j (A_j - A_i) \mathbf{L}_i \nabla W_{ij}
$$
 (1)

with the correction matrix [\[Bonet and Lok](#page-1-1) [1999\]](#page-1-1)

$$
\mathbf{L}_{i} = \left(\sum_{j \in \mathcal{N}_{i}^{0}} V_{j} \nabla W_{ij} \otimes \mathbf{X}_{ji}\right)^{-1}.
$$
 (2)

Authors' addresses: Tassilo Kugelstadt, kugelstadt@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany; Jan Bender, bender@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany; José Antonio Fernández-Fernández, fernandez@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany; Stefan Rhys Jeske, jeske@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany; Fabian Löschner, loeschner@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany; Andreas Longva, longva@cs.rwth-aachen.de, RWTH Aachen University, Aachen, Germany.

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#### <span id="page-1-0"></span>1.1 Computation of Matrix  $D_i$

The matrix  $D_i \in \mathbb{R}^{9 \times 3n}$ , which was introduced in Eq. [\(12\)](#page-0-0) to compute the deformation gradient, is a block matrix which is typically sparse. It has a  $9 \times 3$  block for the particle *i* starting at column 3*i*:

$$
(\mathbf{D}_i)_{0,3i} = -\sum_{j \in \mathcal{N}_i^0} V_j \begin{pmatrix} (\mathbf{L}_i \nabla W_{ij})_1 \mathbf{1} \\ (\mathbf{L}_i \nabla W_{ij})_2 \mathbf{1} \\ (\mathbf{L}_i \nabla W_{ij})_3 \mathbf{1} \end{pmatrix}
$$
(3)

and one  $9 \times 3$  block for each rest-pose neighbor particle *j* starting at column 3*j*:

$$
(\mathbf{D}_i)_{0,3j} = V_j \begin{pmatrix} (\mathbf{L}_i \nabla W_{ij})_1 \mathbf{1} \\ (\mathbf{L}_i \nabla W_{ij})_2 \mathbf{1} \\ (\mathbf{L}_i \nabla W_{ij})_3 \mathbf{1} \end{pmatrix} . \tag{4}
$$

Note that all components of the matrix  $D_i$  only depend on quantities from the rest pose which means that they are constant during the simulation.

### 1.2 Computation of Matrix  $H_{ii}$

In Eq. [\(21\)](#page-0-0) the matrix  $H_{ij} \in \mathbb{R}^{3\times3n}$  was introduced to compute the error vectors  $\mathcal{E}^i_{ij}$ .  $H_{ij}$  is also a block matrix and typically sparse. It has the a  $3 \times 3$  block for particle *i* starting at column 3*i*:

$$
(\mathbf{H}_{ij})_{0,3i} = -\sum_{k \in \mathcal{N}_i^0} V_k \left(\mathbf{L}_i \nabla W_{ik}\right)^T \mathbf{X}_{ij} \mathbb{1} - \mathbb{1}.
$$
 (5)

For the rest-pose neighbor particle  $j = k$  we get

$$
(\mathbf{H}_{ij})_{0,3j} = V_k \left( \mathbf{L}_i \nabla W_{ij} \right)^T \mathbf{X}_{ij} \mathbb{1} + \mathbb{1},\tag{6}
$$

and finally for the neighbor particle  $k \neq j$  we get

$$
(\mathbf{H}_{ij})_{0,3k} = V_k (\mathbf{L}_i \nabla W_{ik})^T \mathbf{X}_{ij} \mathbb{1}.
$$
 (7)

Note that the matrix  $H_{ij}$  is constant during the simulation since its components only depend on quantities from the rest pose.

#### REFERENCES

<span id="page-1-1"></span>J. Bonet and T.-S. L. Lok. 1999. Variational and momentum preservation aspects of Smooth Particle Hydrodynamic formulations. Computer Methods in Applied Mechanics and Engineering 180, 1 (1999), 97 – 115.